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Studies of the Sun with Hinode and Numerical Simulations

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based on Iijima, H., 2016, DSc thesis, UTokyo Iijima & TY, 2015, ApJL, 812, L30

1992/01/12







Hinode since 2006 Solar Optical Telescope (SOT)

- FG (filtergram): High-resolution images of the photosphere and chromosphere
- SP (Spectropolarimeter): high-precision spectropolarimetric diagnostics of the photosphere

EUV Imaging Spectrometer (EIS)

diagnostics of the temperature, velocity, and density in the transition region and the corona

X-Ray Telescope (XRT)

high-resolution images of the corona

The solar atmosphere



(figure from Phillips, Feldman, Landi 2008)



Courtesy T. J. Okamoto, Hinode SOT, JAXA / NAOJ

Alfvenic waves in the chromosphere

Hinode Observations:

– Okamoto+ (2007), Okamoto+, TY (2015)

Numerical Simulations:

- Antolin, TY, Van Doorsselaere (2014), Antolin+, TY (2015)

The solar atmosphere



(figure from Phillips, Feldman, Landi 2008)

Coronal transverse waves in a prominence Okamoto+ 2007



Coronal transverse waves in a prominence

(km)

0

(km)

5,000

4,000

3,000

2,000

1,000

5,000

4,000

3,000

2.000

1,000

0

19:44

0

19:44

С

(S2)

S1

10.000

B (S1)

S2 S3 S4

20,000

S5

30,000

10,000

Okamoto+ 2007

fine-scale threadlike

structures oscillating in the plane of the sky with periods of several minutes

Alfven waves

(c.f. Ofman & Wang 2008)

Fig. 1. High-resolution image on the solar limb obtained with SOT aboard Hinode. This observation was performed with a cadence of 15 s from 19:33 to 20:44 UT on 9 November 2006. Tick marks have a spacing of 1000 km on the Sun. A radial density filter is applied to show the brighter photosphere and the fainter coronal structures in the same image. The main sunspot of NOAA AR 10921 as well as the trailing bright plage areas are vis ible on the disk. Above the limb. ubiquitous vertical spicules are seen below the horizontal threads of the AR prominence. The cloudlike prom inence structure is located 10.000 to 20.000 km above the visible limb and exhibits a very complex



fine structure with predominant horizontal threadlike features. The intensity of the prominence in Ca II H-line radiation is about 1% of the on-disk photosphere.





19:46

19:46

19:48

19:48

(UT)

19:50

19:50

sides of the long thread oscillating synchronously. (B to F) Height-time plots (shown in negative contrast) for the locations indicated in (A). Maximum and minimum amplitudes occur at nearly the same time for all locations.

Fine strand-like structures in the oscillating loop

Antolin, TY, Van Doorsselaere, 2014, ApJ

The transverse oscillations can lead to KH instabilities that deform the cross-section area of the loops. The vortices generated from the instability are velocity sheared regions with enhanced emissivity hosting current sheets.



Resonant absorption of transverse oscillation and associated heating in a prominence

Okamoto+, TY 2015; Antolin+, TY 2015

a Call H









Turbulence in the prominence

Hinode

- Berger+ (2008, 2010)

Numerical simulations

- Hillier+ (2012), Kaneko & TY (2016a,b)

Turbulence in a quiescent prominence Berger+ 2008; 2010

30-Nov-2006 03:30:14 UT

Turbulence in a quiescent prominence

Berger+ 2008; 2010

filamentary downflows and vortices

dark, episodic upflows: 170-700 km in width, exhibit turbulent flow and rise with constant speeds of 20 km/s from the base to heights of 10-20 Mm. resemble buoyant starting plumes



Magnetic Rayleigh-Taylor instability in a prominence

Hillier+ (2012)

Magnetically supported prominence plasma suffers from the RT instability. It may evolve into the turbulence state in the nonlinear phase. Demonstration by MHD simulations.





Magnetic RT instability in a prominence

Kaneko & TY (2016a, b in prep.)

3D MHD simulations based on the newly proposed "reconnection condensation" model. The simulation includes the optically-thin radiative cooling and thermal conduction effects.

The formed prominence shows a turbulent structure probably driven by the magnetic Rayleigh-Taylor instability.



Studies on chromospheric jets

Hinode Observations:

- De Pontieu+ (2007), Pereira+ (2012)

Numerical Simulations:

- Iijima & TY (2015), Iijima (2016, PhD thesis UTokyo)

Chromospheric jets

- Spicules in quiet regions and coronal holes
- Dynamic fibrils in active regions

Manifestation of dynamic plasma processes

- Plasma flows (super-sonic speed)
- Magnetic fields (from high- to low-beta regime)
- MHD waves, Shock waves (mode conversion, non-linear procs.)
- Thermal processes: Radiative cooling, shock heating
- Ionization, recombination ...

Clue for understanding the transport of energies to the corona

Courtesy T. J. Okamoto, Hinode SOT, JAXA / NAOJ





De Pontieu+ (2007)

Statistical studies of chromospheric jets

Lengths, velocities, lifetimes are different among different classes of jets.



AR(hot) dynamic fibrils (De Pontieu+ 2007)



Chromospheric Alfvenic Waves

De Pontieu+ 2007

Alfven waves with amplitudes in the order of 10 to 25 km/s and periods of 100 to 500 sec. energetic enough to accelerate the solar wind and possibly to heat the quiet corona (c.f. Okamoto & De Pontieu 2011)



Chromospheric jets: model



(figure from lijima 2016, PhD UTokyo)

The dense cool chromospheric plasma is lifted by unknown mechanism(s) into the relatively tenuous corona. It is guided by a vertical magnetic field and is observed as an elongated jet.

The rebound shock model for jet driving



Hollweg (1982)

(also Hollweg+ 1982; Sterling & Hollweg 1988; Suematsu+ 1982, Kudoh & Shibata 1999, Takasao+ 2013, etc.)

The transition region (contact discontinuity) is lifted up by interactions with the rebound shock trains propagating in the chromosphere.



The shock trains can be generated by various processes: from convective overshoot, Alfven waves, reconnection events etc.

Chromospheric jets: model



(figure from lijima 2016, PhD UTokyo)

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Radiative MHD simulations of chromospheric jets lijima & TY (2015); lijima (2016, PhD UTokyo)

Typical case: $T_c = 0.4$ MK, $B_0 = 30$ G



Radiative MHD simulations of chromospheric jets lijima & TY (2015); lijima (2016, PhD UTokyo)

Typical case: $T_c = 0.4$ MK, $B_0 = 30$ G



Rebound-shock ejection of jets lijima & TY (2015); lijima (2016, PhD UTokyo)

1D distribution along a field line passing through the top of a jet







lijima & TY (2016); lijima (2016, PhD UTokyo)



magnetic field lines & T color

stream lines & |V| color



Hinode results on chromospheric dynamics

Waves

Transverse waves found in prominences (Okamoto+ 2007), spicules (De Pontieu+ 2007), evidence of the resonant-absorption thermalization (Okamoto+ 2007)

Prominence internal flows

Turbulent flows inside (Berger+ 2008; 2010)

Chromospheric jets (Spicules)

Extensive statistical studies are carried out.

 \rightarrow By combination with radiative MHD simulations, understanding of the driving mechanisms are going to be achieved.

Chromospheric reconnection

Anemone jets (Shibata+ 2007), penumbral micro-jets (Katsukawa+ 2007)

 \rightarrow Ubiquitous magnetic reconnection

End

Basic equations

lijima & TY (2015); lijima (2016, PhD UTokyo)

Magnetohydrodynamic equations

$$\begin{split} \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{V}) &= 0\\ \frac{\partial (\rho \mathbf{V})}{\partial t} + \nabla \cdot \left[\rho \mathbf{V} \otimes \mathbf{V} + \left(p + \frac{B^2}{8\pi} \right) \mathbf{I} - \frac{\mathbf{B} \otimes \mathbf{B}}{4\pi} \right] &= \rho \mathbf{g}\\ \frac{\partial e_{\text{tot}}}{\partial t} + \nabla \cdot \left[\left(e_{\text{tot}} + p + \frac{B^2}{8\pi} \right) \mathbf{V} - \frac{1}{4\pi} \mathbf{B} \left(\mathbf{V} \cdot \mathbf{B} \right) \right] \\ &= \rho \left(\mathbf{g} \cdot \mathbf{V} \right) + Q_{\text{cnd}} + Q_{\text{rad}}\\ \frac{\partial \mathbf{B}}{\partial t} + \nabla \cdot \left(\mathbf{V} \otimes \mathbf{B} - \mathbf{B} \otimes \mathbf{V} \right) = 0 \end{split}$$

EoS for LTE plasma (tabulated)

Radiative cooling lijima & TY (2015); lijima (2016, PhD UTokyo)

Total cooling term is switched according to the vertical column density

 $Q_{\rm rad} = \left[1 - \xi\left(m_c\right)\right] Q_{\rm thick} + \xi\left(m_c\right) Q_{\rm thin}$

$$Q_{\text{thick}} = e^{-(\tau/\tau_0)^2} Q_{\text{J}} + \left[1 - e^{-(\tau/\tau_0)^2}\right] Q_{\text{F}}.$$
$$Q_{\text{F}} = -\nabla \cdot \boldsymbol{F}, \quad \boldsymbol{F} = \int_{4\pi} \boldsymbol{n} I(\boldsymbol{n}) d\boldsymbol{n}$$
$$Q_{\text{J}} = 4\pi \alpha_{\text{R}} \left(J - S\right), \quad J = \frac{1}{4\pi} \int_{4\pi} I(\boldsymbol{n}) d\boldsymbol{n}.$$

Intensity *I* (w/ gray approx.) is obtained by solving radiative transfer eq.

$$\frac{dI}{ds} = \alpha \left(B - I \right)$$

$$Q_{\rm thin} = -n_H n_e \Lambda(T)$$

radiative loss function: CHIANTI database with extension by Goodman & Judge (2012) for low temperature plasma



20:00:37 UT



Found in active regions Cusp-shaped structure and bright footpoint length 1-4 Mm, lifetime 100-500s velocity 5-20km/s ~local Alfven speed

Coronal X-ray Jets found in early 1000's



Shibata et al. 1992

TY & Shibata (1995, 1996)





Waves as a carrier of energy



The solar atmosphere



Wave as a carrier of the energy

